

Introduction to Photolysis  
AOSC 433/633 & CHEM 433

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**Class Web Site:** <http://www.atmos.umd.edu/~rjs/class/spr2015>

**Lecture 10**  
**10 March 2015**

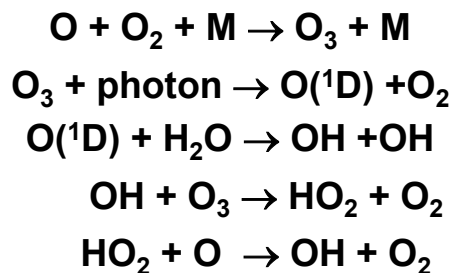
# Importance of Radicals

- With a few exceptions, the only reactions between molecules that proceed at appreciable rates are those involving at least one radical
- Radicals require significant energy to form: a bond must be broken
- Radical formation is tied to absorption of photons that “photodissociate” a compound, leading to radical formation

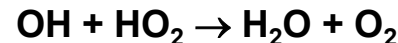
## Initiation



## Propagation

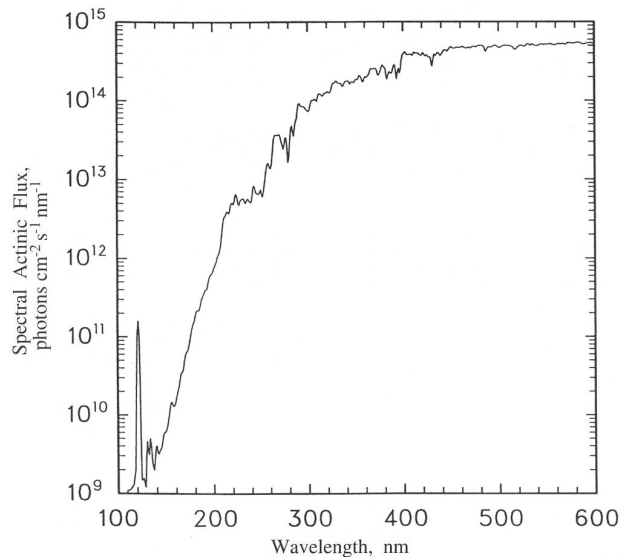


## Termination



- Motivation for Today's Lecture:

**a) How does atmosphere go from this:**



From DeMore et al., *Chemical Kinetics and Photochemical Data for Use in Stratospheric Modeling*, Evaluation No. 11, 1994.

**to this ?**

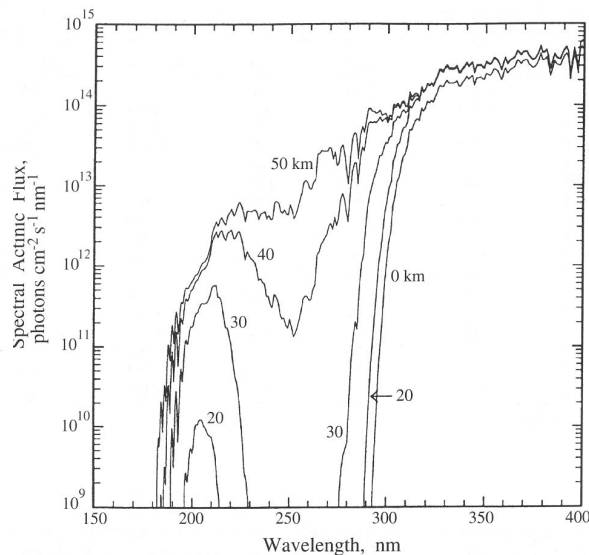
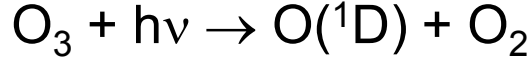


FIGURE 3.3 Solar spectral actinic flux (photons cm<sup>-2</sup> s<sup>-1</sup> nm<sup>-1</sup>) at various altitudes and at the Earth's surface (DeMore et al., 1994).

From Seinfeld and Pandis, *Atmospheric Chemistry and Physics*, 1998.

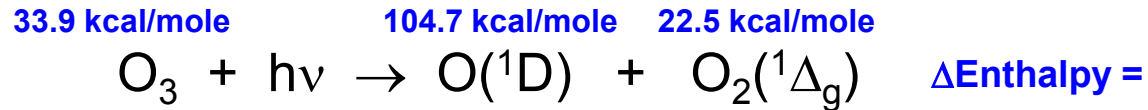
**b) Biological Effects of UV Radiation**

# Energetics of Photolysis



$h\nu$  represents a photon with specific energy.

Let's examine enthalpy of this reaction:



Photon Energy:

$$\varepsilon = \frac{hc}{\lambda} \Rightarrow \lambda_{\text{max}} = \frac{hc}{\Delta\text{Enthalpy}}$$

For  $\text{O}_3$  photo-dissociating to  $\text{O}(^1\text{D})$ :

$$\lambda_{\text{max}} = \frac{hc}{\Delta\text{Enthalpy}} = \frac{2.85 \times 10^4 \text{ kcal/mole nm}}{\Delta\text{Enthalpy}} = \text{-----} =$$

# Energetics of Photodissociation



**Atomic oxygen:** (Note: you will not be “responsible” for the material below on any exam ☺)

**Ground state** – two unpaired electrons in the 2p orbitals:  $(1s)^2(2s)^2(2p_1)^2(2p_2)^1(2p_3)^1$

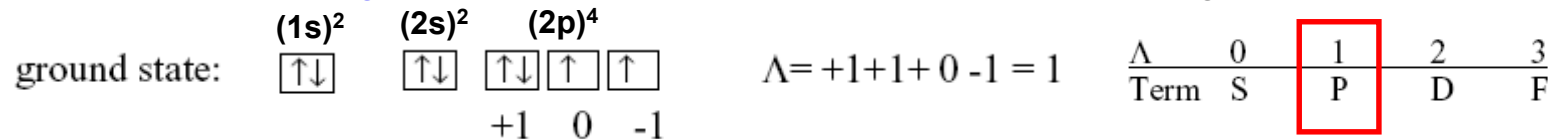
Called  $^3\text{P}$ :

“3” represents  $2S+1$ , where S is spin of all of the unpaired electrons.

There are 2 unpaired electrons, each with spin of  $\frac{1}{2}$

Hence,  $S = 1$  and  $2S+1 = 3 \leftarrow$  spin angular momentum

**P** represents orbital angular momentum, found from an electron diagram of filled orbitals:



**Excited state** – one electron moves from  $2p_3$  to  $2p_2$  :  $(1s)^2(2s)^2(2p_1)^2(2p_2)^2$

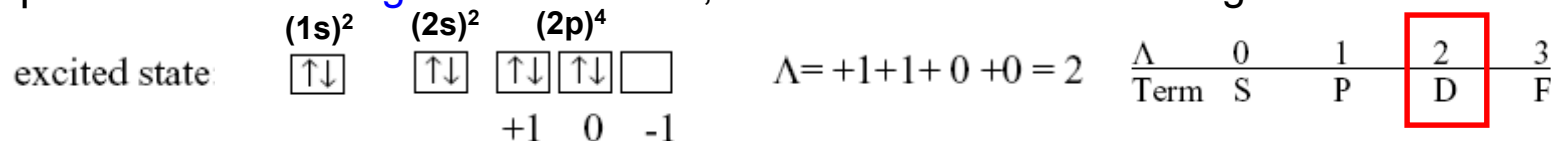
Called  $^1\text{D}$ :

“1” represents  $2S+1$ , where S is spin of all of the unpaired electrons.

There are no unpaired electrons!

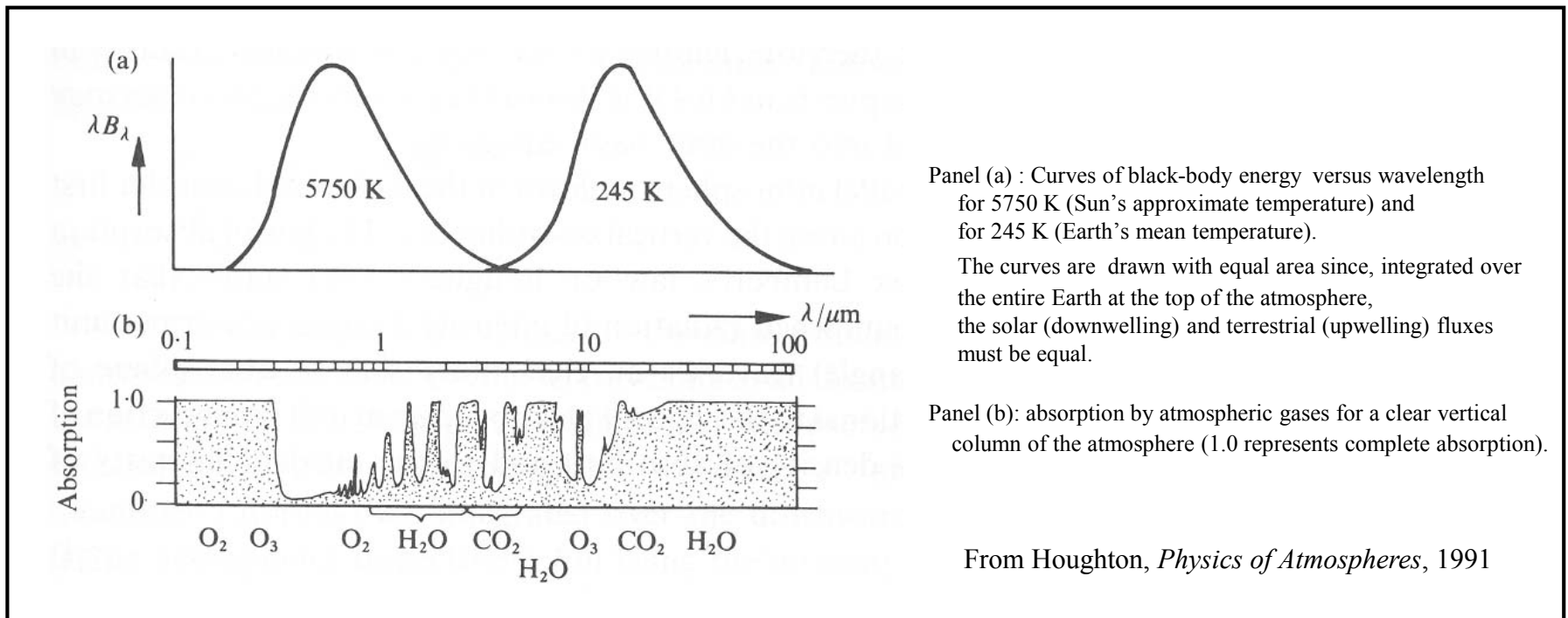
Hence,  $S = 0$  and  $2S+1 = 1 \leftarrow$  spin angular momentum

**D** represents orbital angular momentum, found from an electron diagram of filled orbitals:



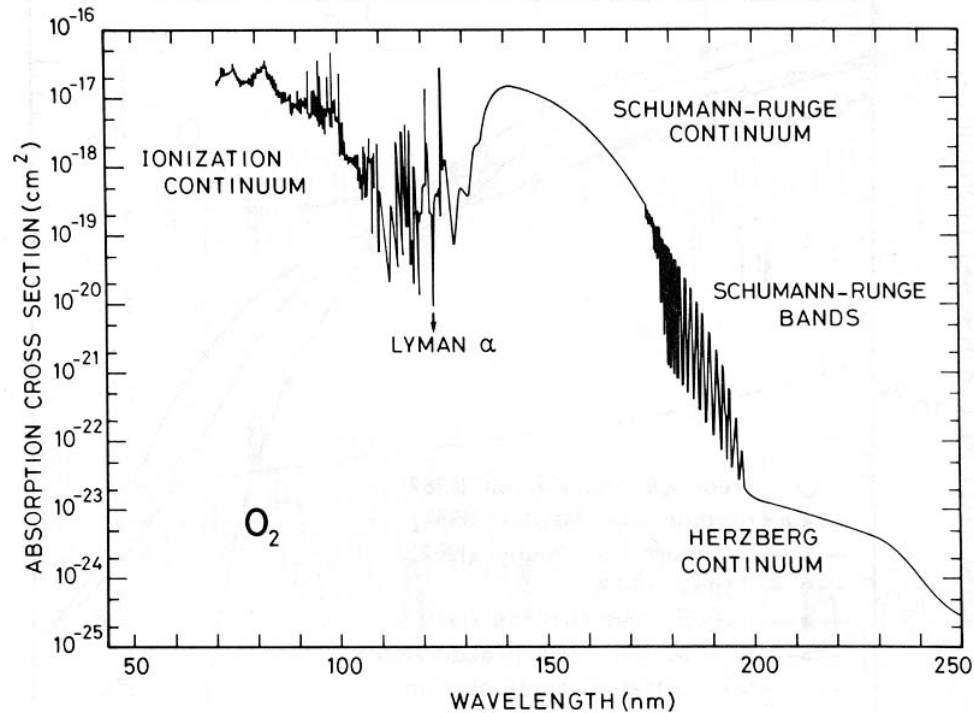
# Atmospheric Radiation

- Solar irradiance (downwelling) at top of atmosphere occurs at wavelengths between  $\sim 200$  and  $2000$  nm ( $\sim 5750$  K “black body” temperature)



- Absorption and photodissociation in the UV occurs due to changes in the electronic state (orbital configuration) of molecules

# Absorption Cross Section of O<sub>2</sub>



From Brasseur & Solomon, *Aeronomy of the Middle Atmosphere*, 1986

- O<sub>2</sub> can not dissociate longward of ~250 nm
- All absorption shown above is dissociative (e.g., leads to production of two O atoms)
- Structure in the O<sub>2</sub> cross section is related to whether the initial transition involves an unbound electronic state (smooth) or involves a specific vibrational level of an electronic state (banded, due to requirement of specific quanta of energy)

# Beer-Lambert Law

$$F(z, \lambda) = F_{\text{TOA}}(\lambda) e^{-\tau(z, \lambda)} \quad (\text{TOA : Top of Atmosphere})$$

where:

$$\tau(z, \lambda) = m \int_z^{\infty} \sigma_{\lambda} [C] dz' \quad (\tau: \text{optical depth})$$

F : solar irradiance (photons/cm<sup>2</sup>/sec)

$\sigma_{\lambda}$  : absorption cross section

C : concentration of absorbing gas (molecules/cm<sup>3</sup>)

m : ratio of slant path to vertical path, equal to  $1/\cos(\theta)$  for  $\theta < \sim 75^{\circ}$

$\theta$  : solar zenith angle

Governs basics of radiative transfer in the UV and near IR regions



# Photolysis Frequency

For a specific spectral interval, the photolysis frequency (*partial J value*) of a gas is given by the product of its absorption cross section and the solar irradiance:

$$J_{\text{gas}}(z, \lambda) = \text{Quantum\_Yield}(\lambda) \sigma_{\text{gas}}(\lambda, T) F(z, \lambda)$$

Units:  $\text{s}^{-1} \text{ nm}^{-1}$

The total *photolysis frequency* (*J value*) is found by integrating  $J_{\text{gas}}(z, \lambda)$  over all wavelengths for which the gas photodissociates:

$$J_{\text{gas}}(z) = \int_{\lambda_{\text{min}}}^{\lambda_{\text{max}}} J_{\text{gas}}(z, \lambda) d\lambda$$

Units:  $\text{s}^{-1}$

$$\text{Rate of Reaction} = \frac{d\text{O}_3}{dt} = J [\text{O}_3]; \text{ Units of } J \text{ are } \text{s}^{-1}$$

*More precisely, calculations of photolysis frequencies consider the “spectral actinic flux”, which represents the amount of available photons integrated over all angles, rather than “solar irradiance”. These two quantities differ because of scattering of solar radiation by gases and aerosols, and reflection of radiation by clouds and the surface.*

# Optical Depth of O<sub>2</sub> Absorption

Recall the *Beer-Lambert Law*:

$$F(z, \lambda) = F_{\text{TOA}}(\lambda) e^{-\tau(z, \lambda)} \quad (\text{TOA : Top of Atmosphere})$$

where:

$$\tau(z, \lambda) = m \int_z^{\infty} \sigma_{\lambda} [C] dz' \quad (\tau: \text{optical depth})$$

Also:

$$\int_0^{\infty} [\text{O}_2] dz' \approx 4 \times 10^{24} \text{ molecules/cm}^2$$

O <sub>2</sub> Optical Depth for $\theta = 0^\circ$ , $z = 0$ km		
	$\sigma_{\text{max}} \text{ (cm}^2\text{)}$	$\tau (0 \text{ km})$
Schumann-Runge Continuum		$e^{-\tau (0 \text{ km})}$
Schumann-Runge Bands		
Herzberg Continuum		

# Photolysis Frequency of O<sub>2</sub>

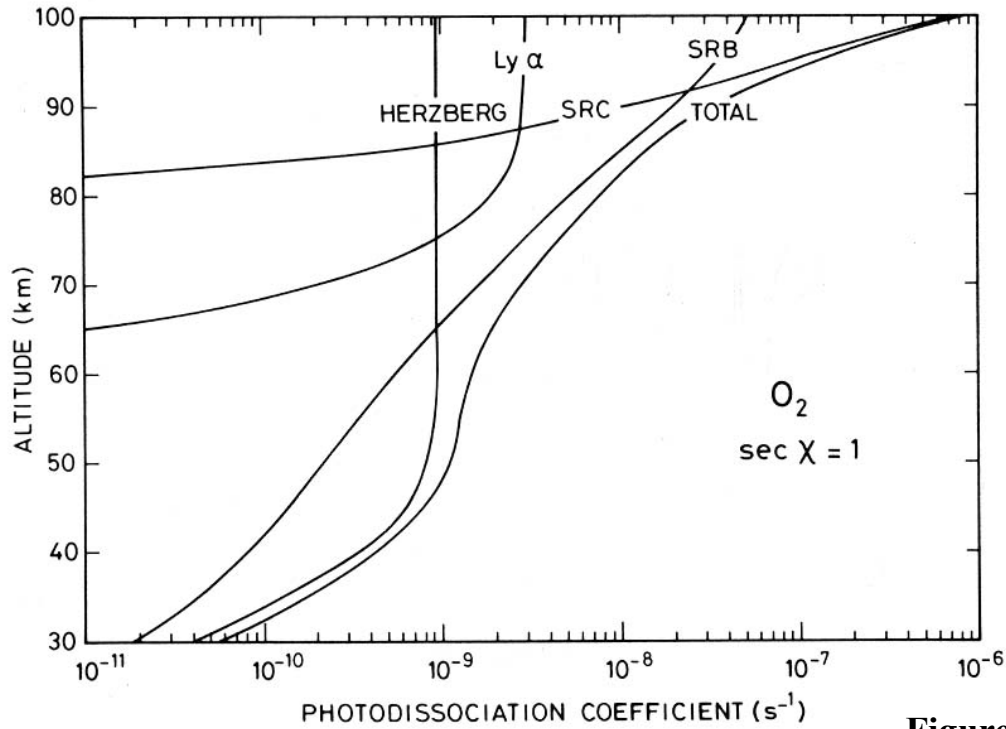


Figure 4.31, Brasseur and Solomon

# Where Does Optical Depth = 1.0 for O<sub>2</sub> ?

$$\tau(z, \lambda) = m \int_z^{\infty} \sigma_{\lambda} [\text{O}_2] dz'$$
$$\approx \sigma_{\lambda} m 4 \times 10^{24} e^{-z/H}$$

Setting  $\tau = 1$  and re-arranging gives:

$$z = H \ln (\sigma_{\lambda} \cdot m \cdot 4 \times 10^{24})$$

Altitude where $\tau = 1$ (for $\theta = 0^\circ$ )		
	$\sigma_{\max}$ (cm <sup>2</sup> )	z ( km)
Schumann-Runge Continuum	$10^{-17}$	
Schumann-Runge Bands	$10^{-20}$	
	$3 \times 10^{-23}$	
Herzberg Continuum	$10^{-23}$	

# Absorption Cross Section of O<sub>3</sub>

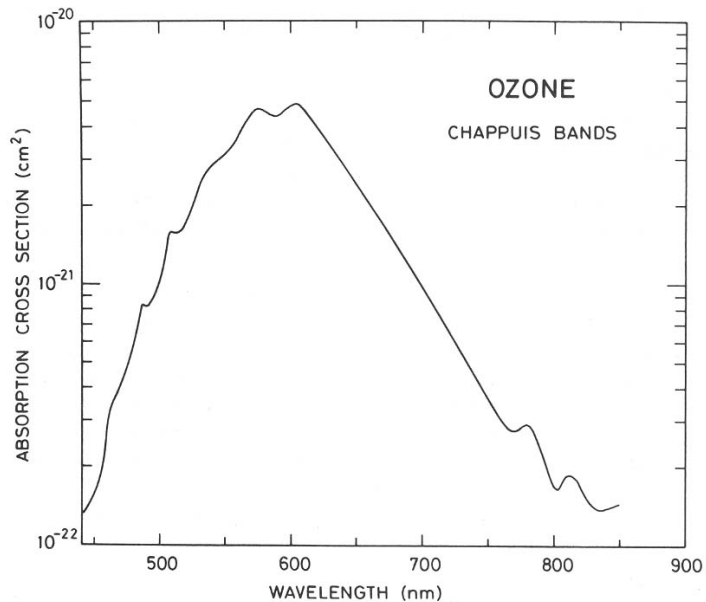
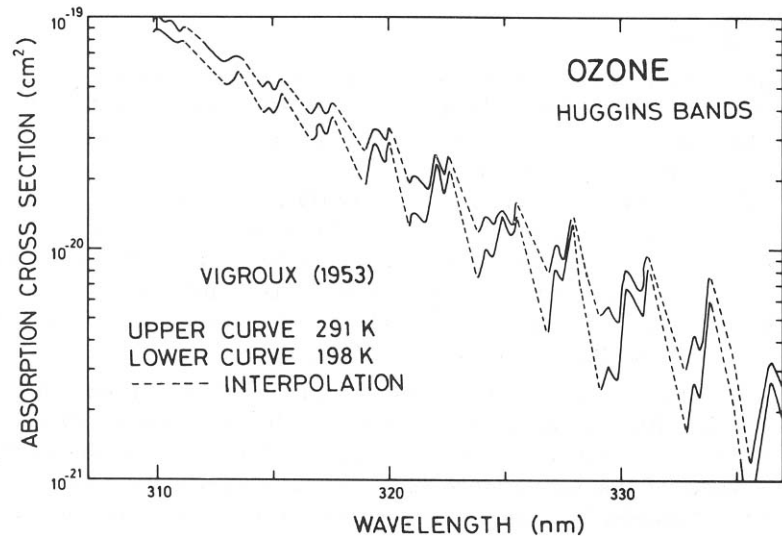
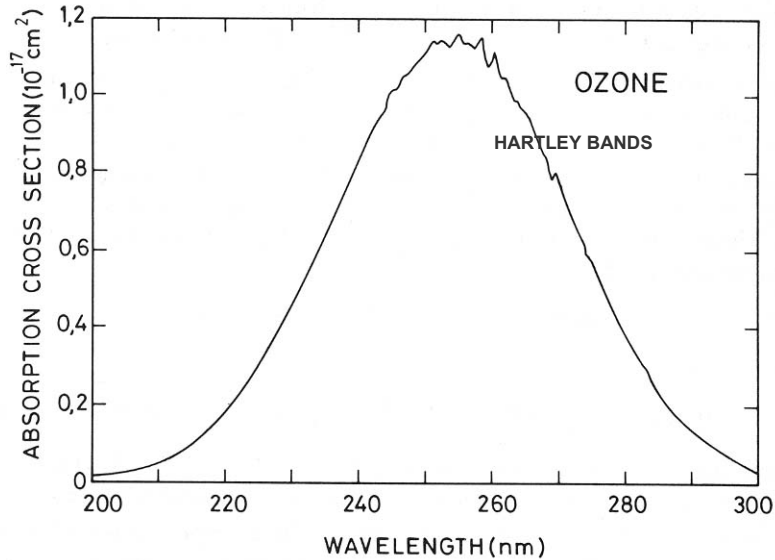


Table 4.6 Theoretical limits corresponding to different photolysis products (nm).

	O <sub>2</sub> ( <sup>3</sup> Σ <sub>g</sub> )	O <sub>2</sub> ( <sup>1</sup> Δ <sub>g</sub> )	O <sub>2</sub> ( <sup>1</sup> Σ <sub>g</sub> +)	O <sub>2</sub> ( <sup>3</sup> Σ <sub>u</sub> +)	O <sub>2</sub> ( <sup>3</sup> Σ <sub>u</sub> -)
O( <sup>3</sup> P)	1180	590	460	230	170
O( <sup>1</sup> D)	410	310	260	167	150
O( <sup>1</sup> S)	234	196	179	129	108

From Brasseur & Solomon, *Aeronomy of the Middle Atmosphere*, 1986

# Optical Depth of O<sub>3</sub> Absorption

A typical mid-latitude column abundance for O<sub>3</sub> is 300 Dobson units (DU):

$$1 \text{ DU} = 2.687 \times 10^{16} \text{ molecules/cm}^2; \quad 300 \text{ DU} = 8 \times 10^{18} \text{ molecules/cm}^2$$

Aside: 
$$\frac{\text{Column O}_3}{\text{Column Air}} = \frac{8 \times 10^{18}}{2 \times 10^{25}} = 0.4 \text{ parts per million} \Rightarrow \text{Ozone is a trace species!}$$

O <sub>3</sub> Optical Depth for $\theta = 0^\circ$ , $z = 0$ km				
	$\sigma_{\text{max}}$ (cm <sup>2</sup> )	$\tau$ (0 km)	$e^{-\tau}$ (0 km)	O <sub>3</sub> Column, $\tau = 1.0$
Hartley (~220 to 280 nm)				
Huggins (~310 to 330 nm)				
Chappuis (~500 to 700 nm)				

# Solar Spectral Actinic Flux

130 ATMOSPHERIC PHOTOCHEMISTRY AND CHEMICAL KINETICS

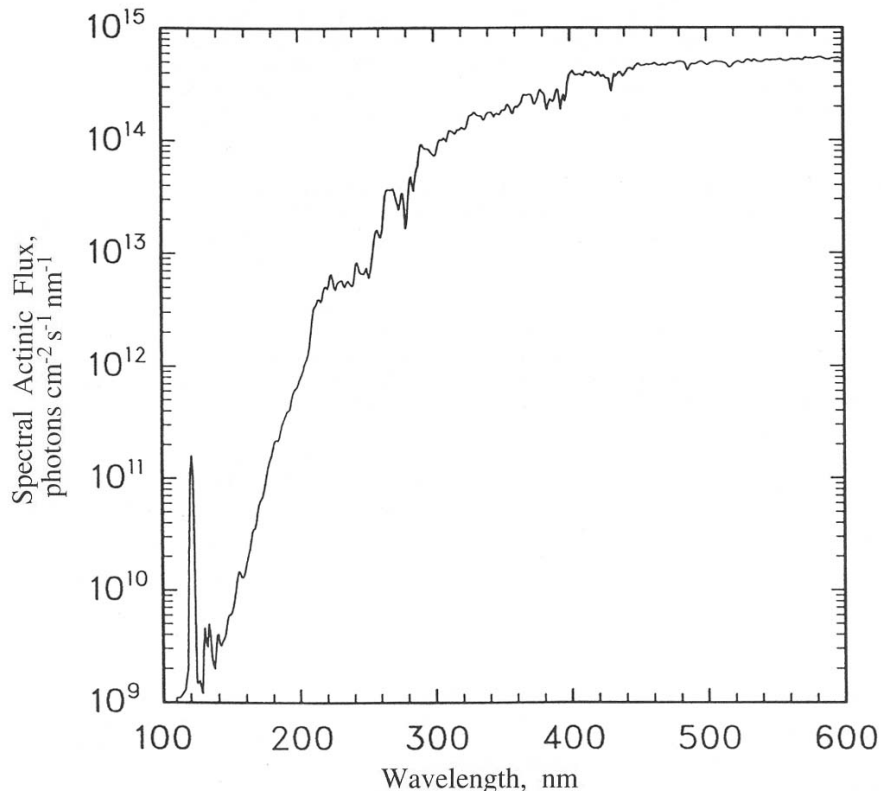


FIGURE 6. Solar spectral actinic flux (photons  $\text{cm}^{-2} \text{s}^{-1} \text{nm}^{-1}$ ) at the top of Earth's atmosphere.

From DeMore et al., *Chemical Kinetics and Photochemical Data for Use in Stratospheric Modeling*, Evaluation No. 11, 1994.

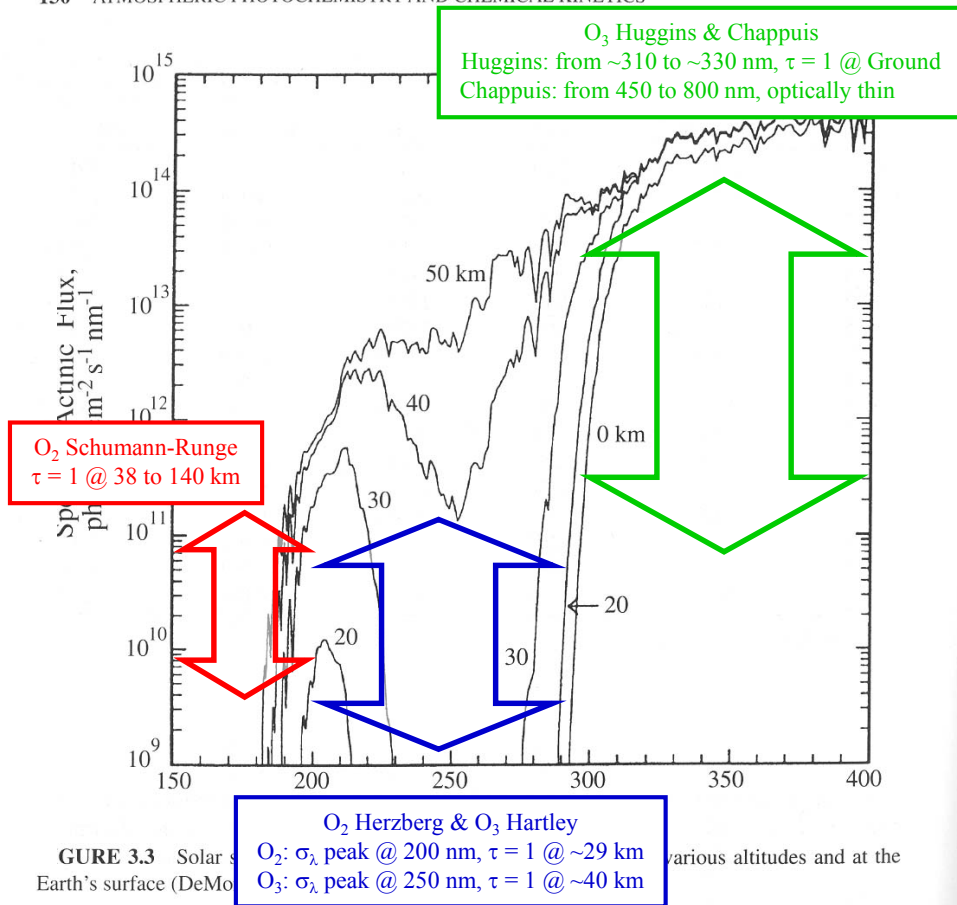
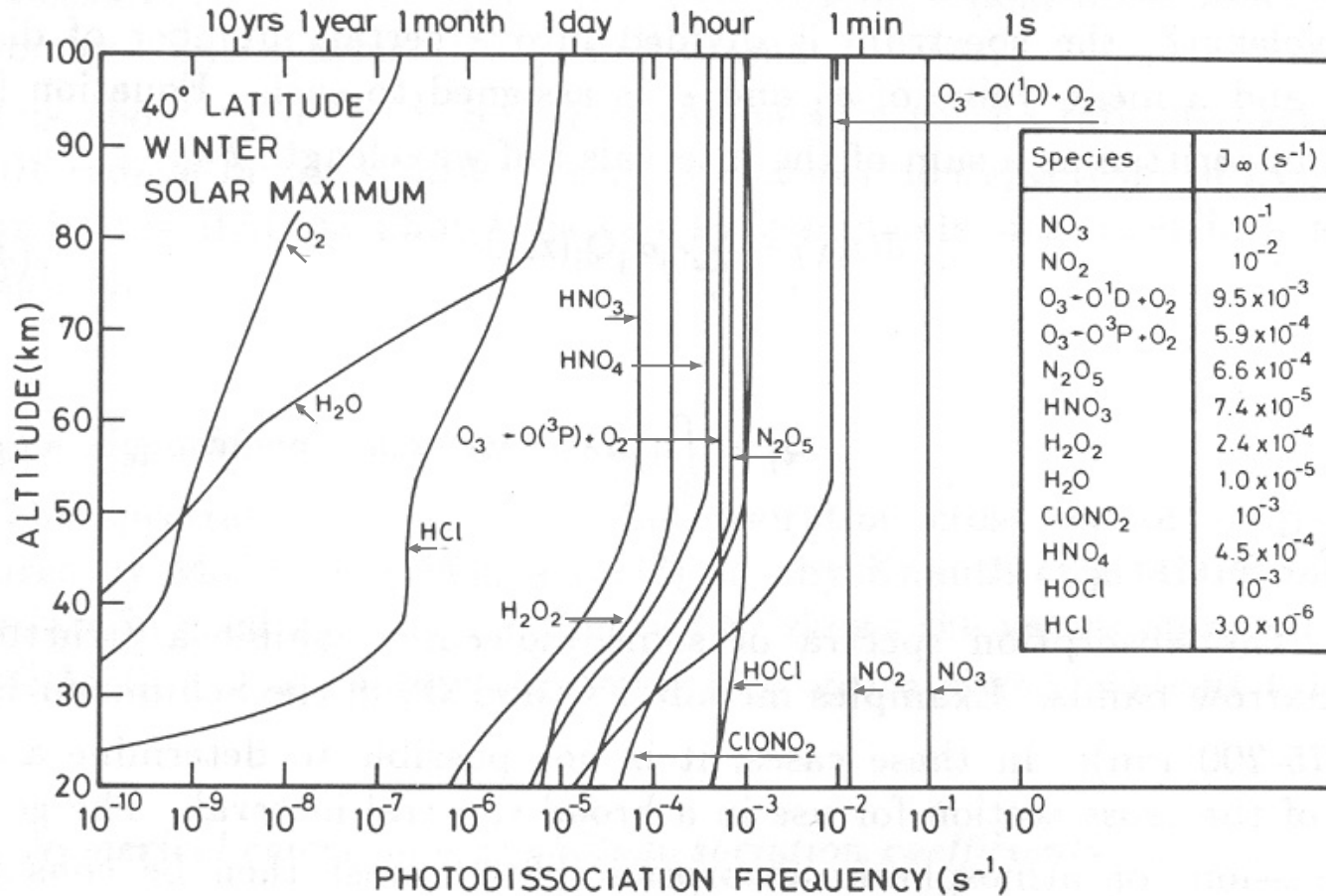


FIGURE 3.3 Solar spectral actinic flux at various altitudes and at the Earth's surface (DeMore et al., 1994).

From Seinfeld and Pandis, *Atmospheric Chemistry and Physics*, 1998.

# Photodissociation Frequencies

**Next goal is to understand:**



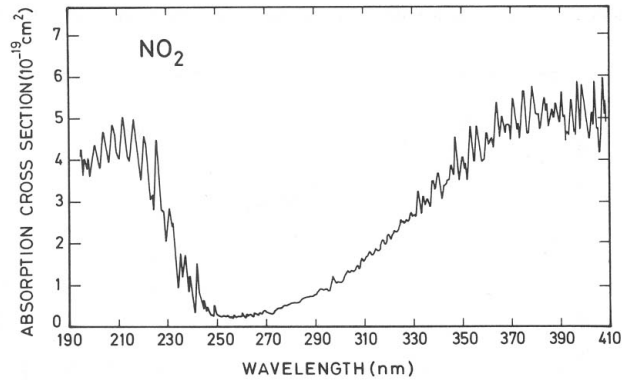
**FIGURE 4.58** Photodissociation frequencies for numerous important atmospheric species.

From Brasseur & Solomon, *Aeronomy of the Middle Atmosphere*, 1986

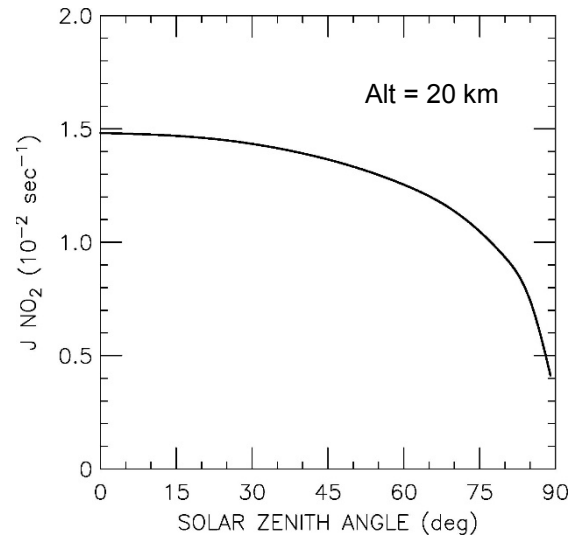
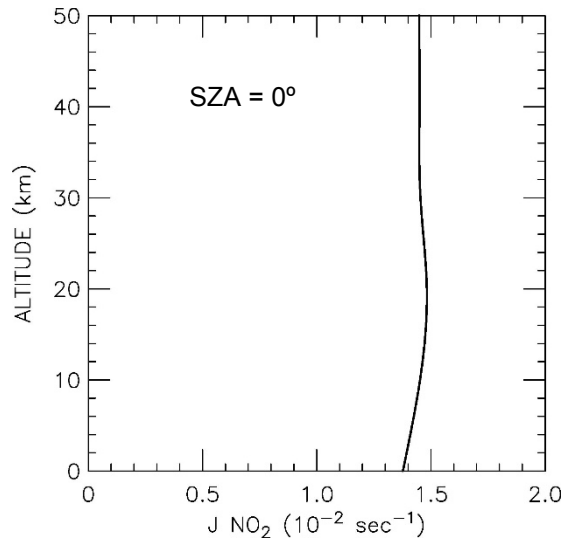


# NO<sub>2</sub> Photolysis

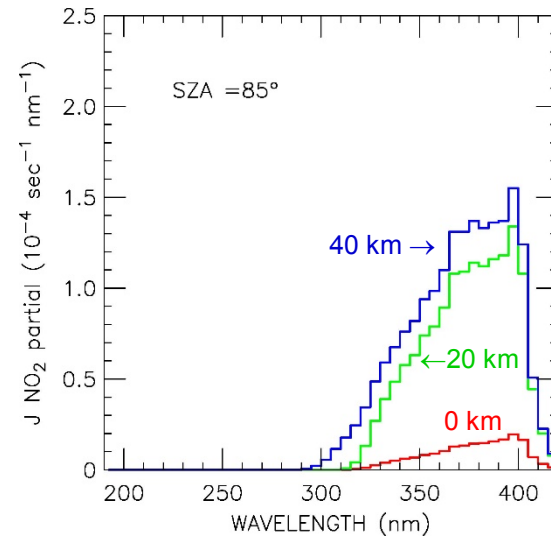
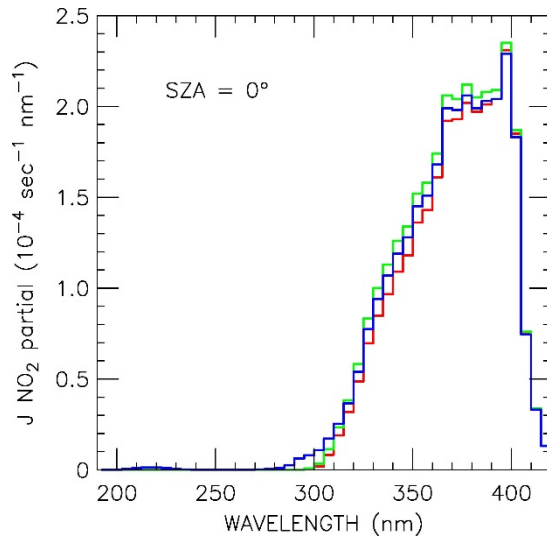
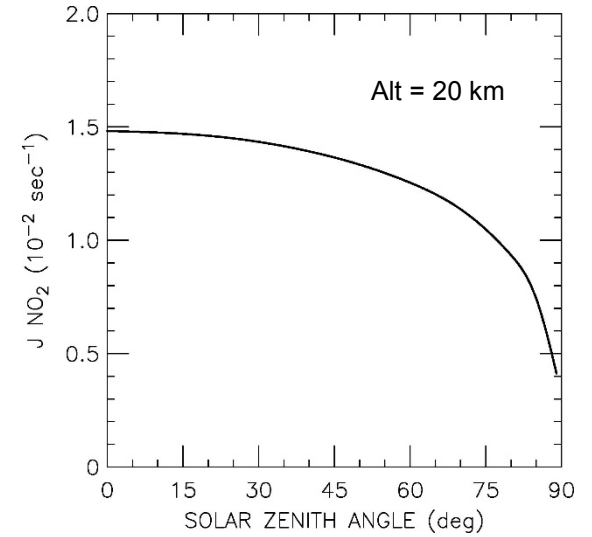
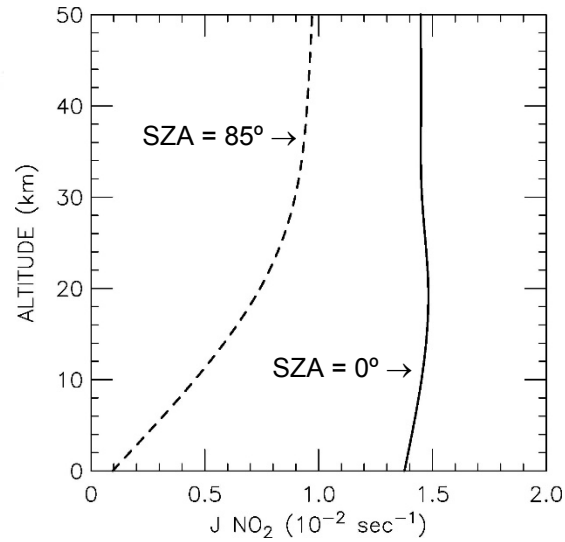
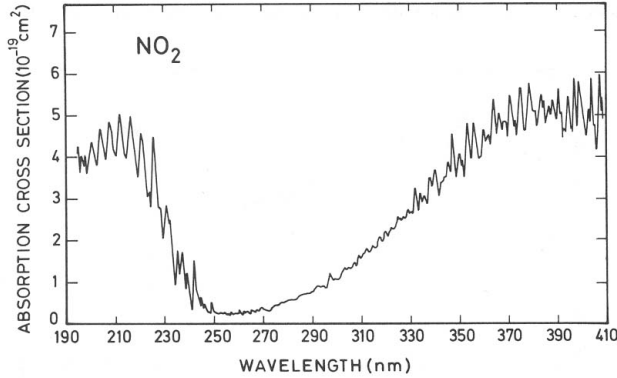
The majority of NO<sub>2</sub> photolysis occurs longward of 300 nm, where the atmosphere is optically thin with respect to absorption by O<sub>3</sub> and O<sub>2</sub>:



leading to a value for  $J_{\text{NO}_2}$  that is nearly independent of height and SZA:

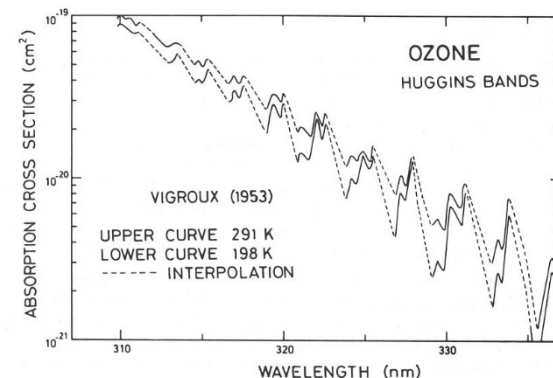
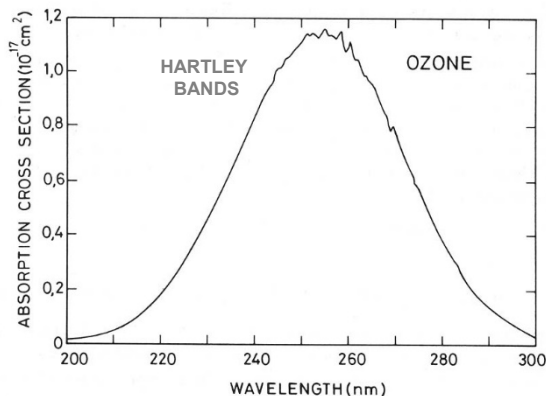
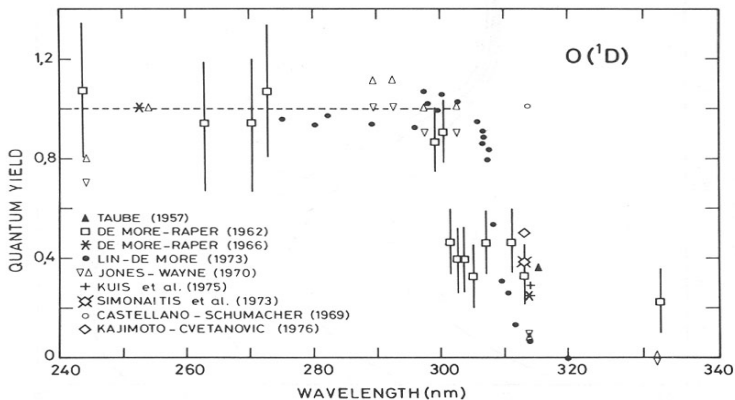


# NO<sub>2</sub> Photolysis

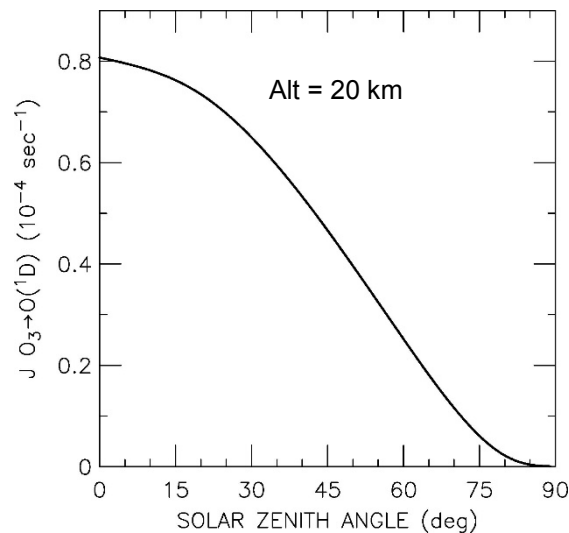
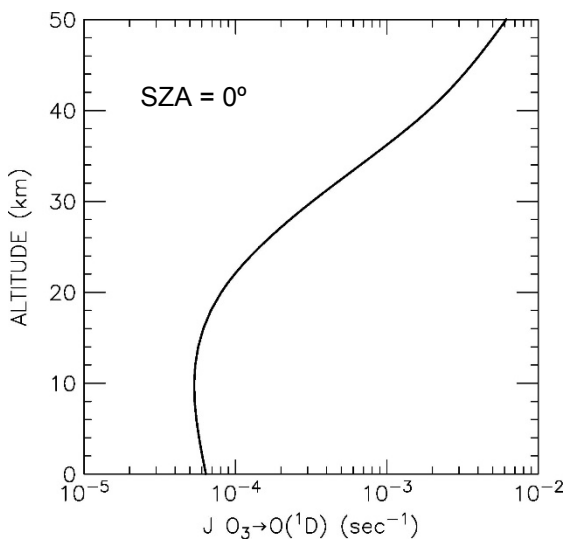


# O<sub>3</sub> → O(<sup>1</sup>D) Photolysis

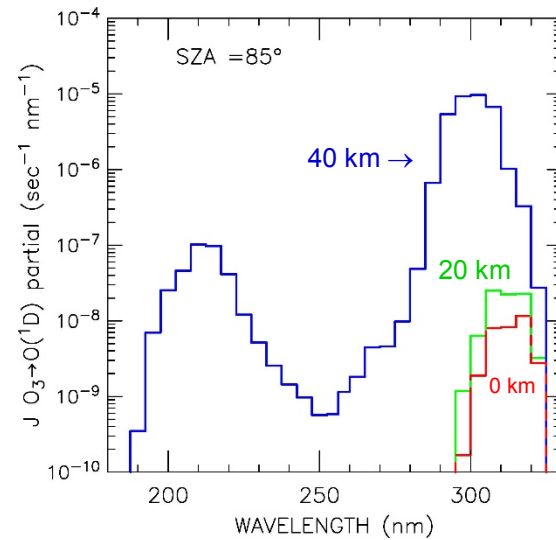
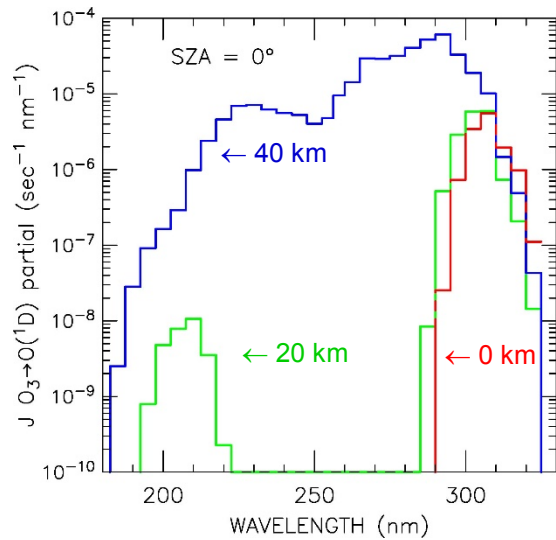
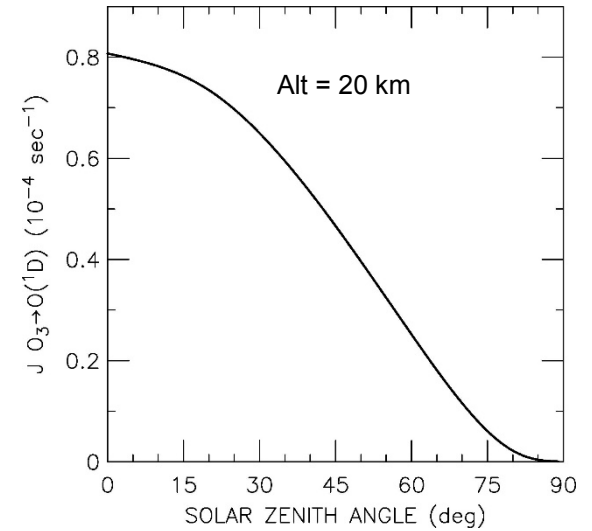
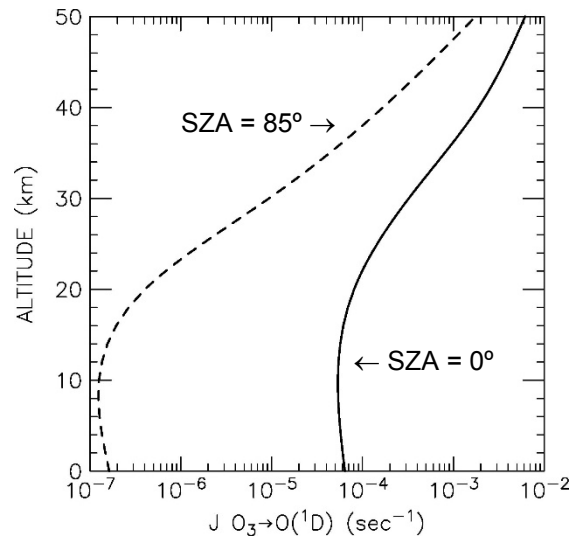
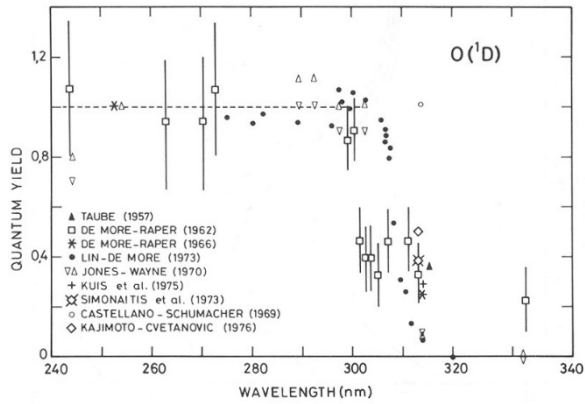
The production of O(<sup>1</sup>D) from photolysis of O<sub>3</sub> occurs shortward of 320 nm, where the atmosphere is basically optically thick with respect to absorption by O<sub>3</sub>:



leading to a value for  $J_{O_3 \rightarrow O(1D)}$  that is dependent on height and SZA:

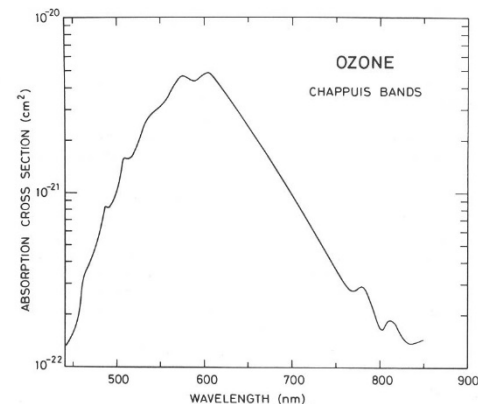
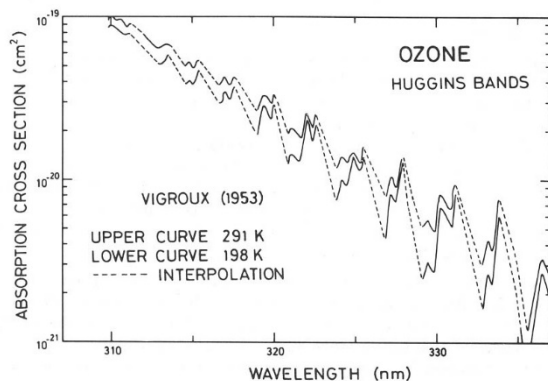
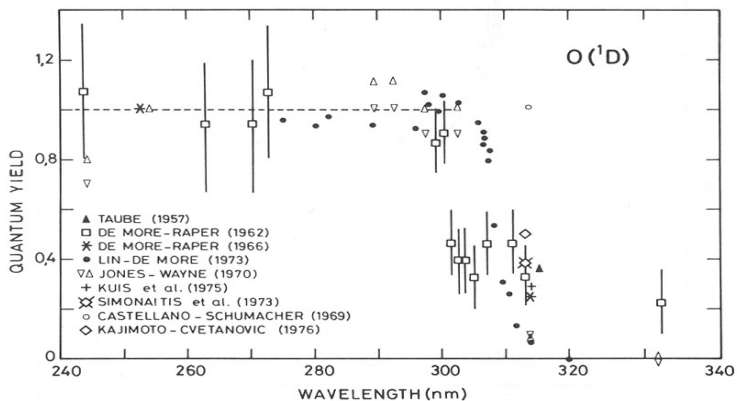


# O<sub>3</sub> → O(<sup>1</sup>D) Photolysis

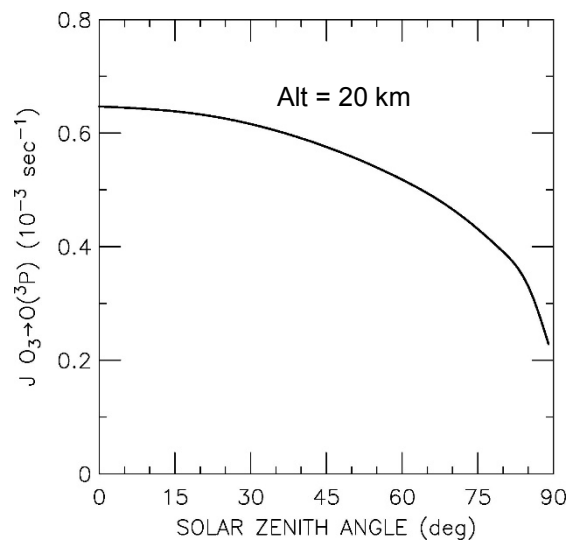
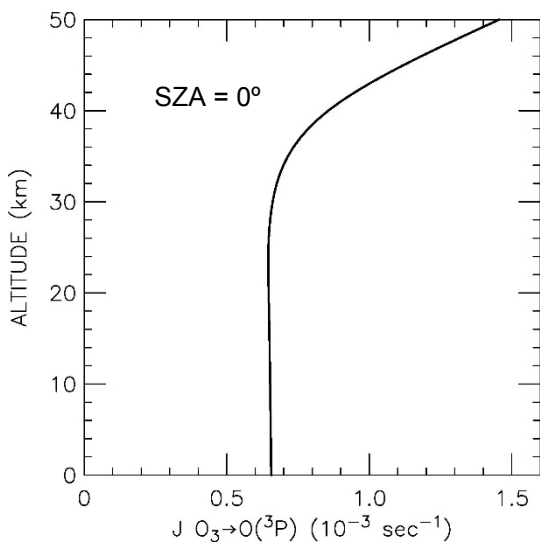


# O<sub>3</sub> → O(<sup>3</sup>P) Photolysis

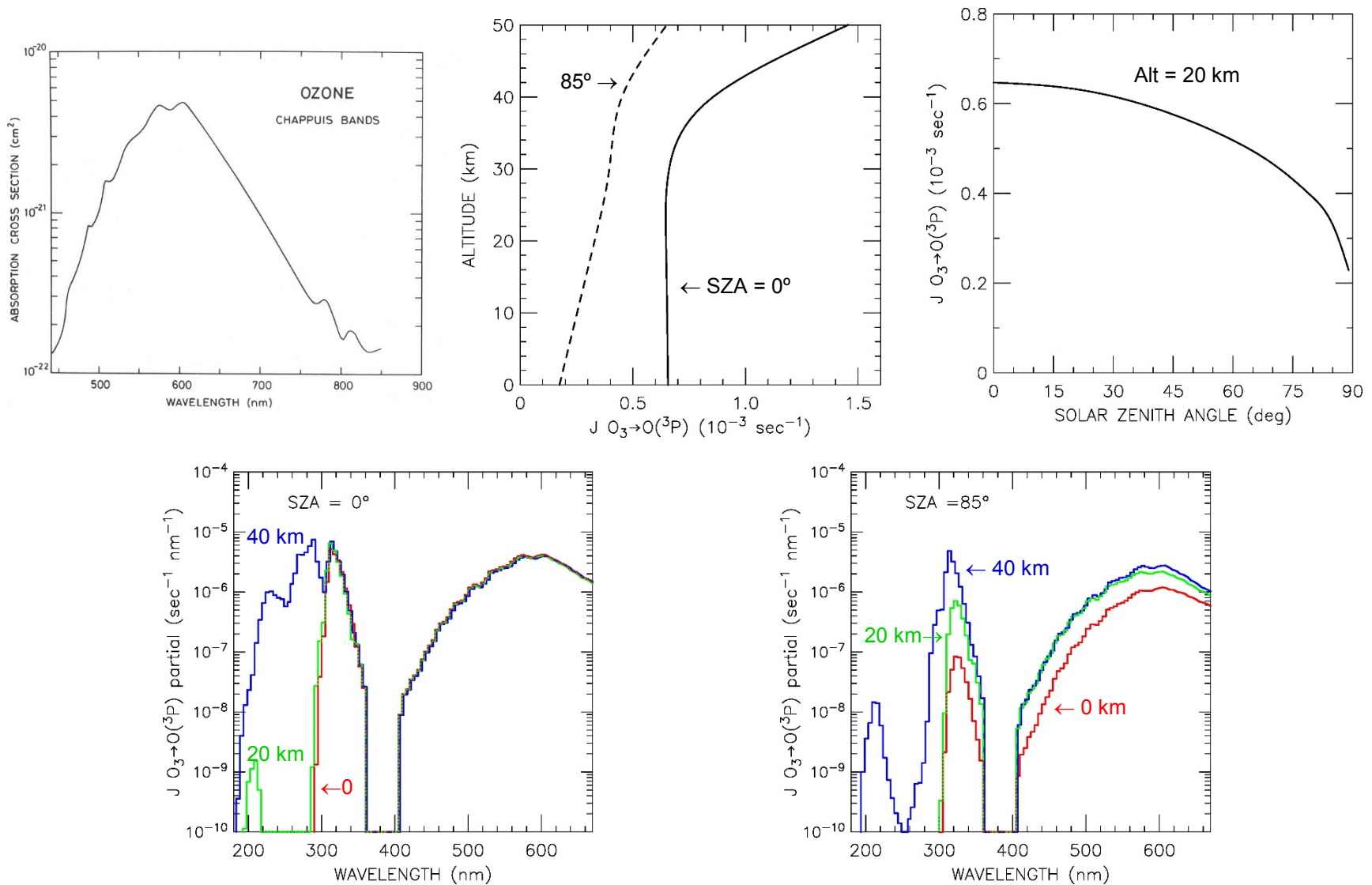
The production of O(<sup>3</sup>P) from photolysis of O<sub>3</sub> occurs mainly longward of 500 nm, where the atmosphere is optically thin with respect to absorption by O<sub>3</sub>:



leading to a value for  $J_{O_3 \rightarrow O(^3P)}$  that is essentially independent of height and SZA:



# $O_3 \rightarrow O(^3P)$ Photolysis



# Biological Effects of UV Radiation

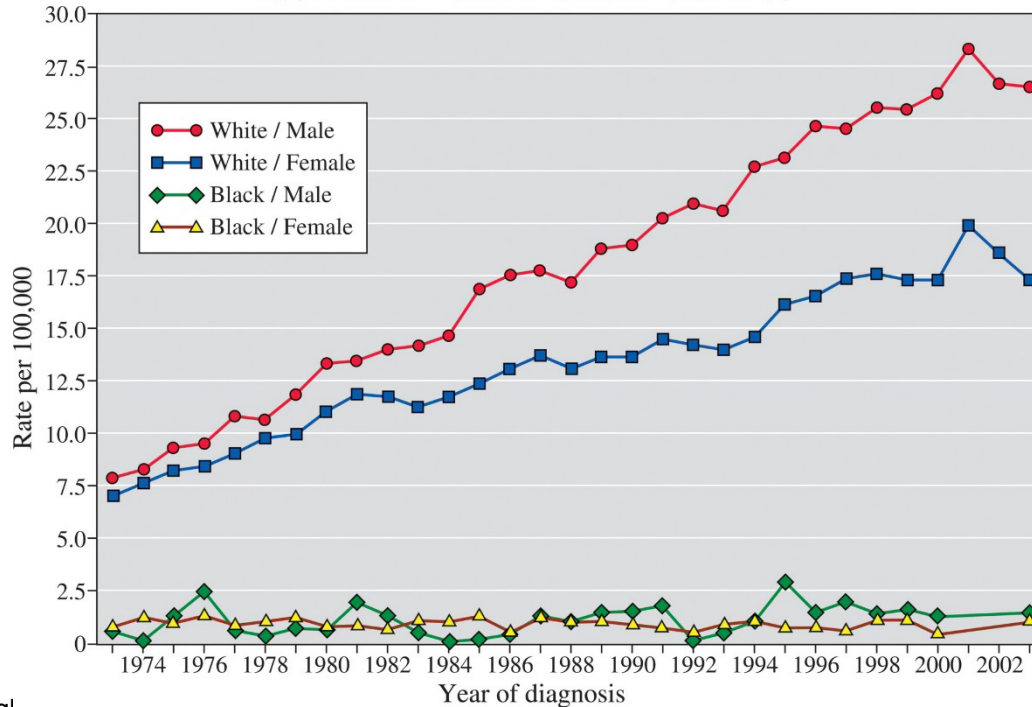
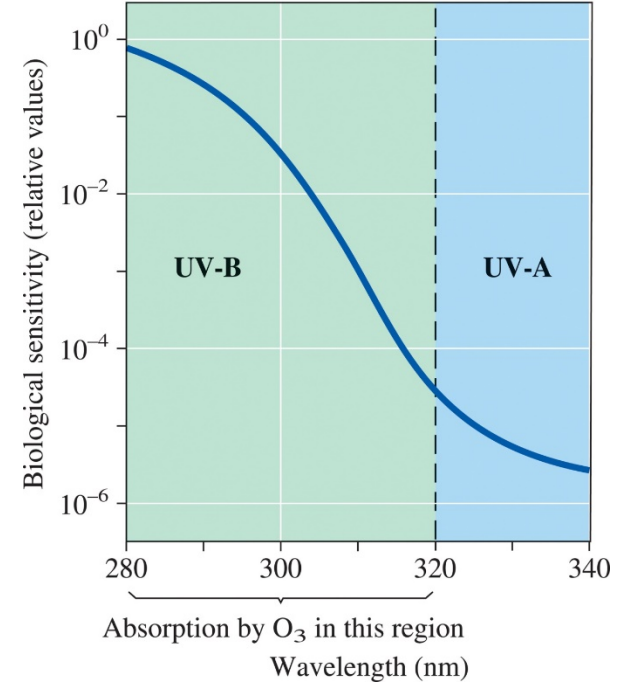
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Table 2.4 Types of UV Radiation			
Type	Wavelength	Relative Energy	Comments
UV-A	320–400 nm	Lowest energy	Least damaging and reaches the Earth's surface in greatest amount
UV-B	280–320 nm	Higher energy than UV-A but less energetic than UV-C	More damaging than UV-A but less damaging than UV-C. Most UV-B is absorbed by O <sub>3</sub> in the stratosphere
UV-C	200–280 nm	Highest energy	Most damaging but not a problem because it is totally absorbed by O <sub>2</sub> and O <sub>3</sub> in the stratosphere

## Chemistry in Context

Figure 2.11, Chemistry in Context

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Increase in incidence of melanoma skin cancer, U.S.  
Figure 2.12, Chemistry in Context

# Biological Effects of UV Radiation

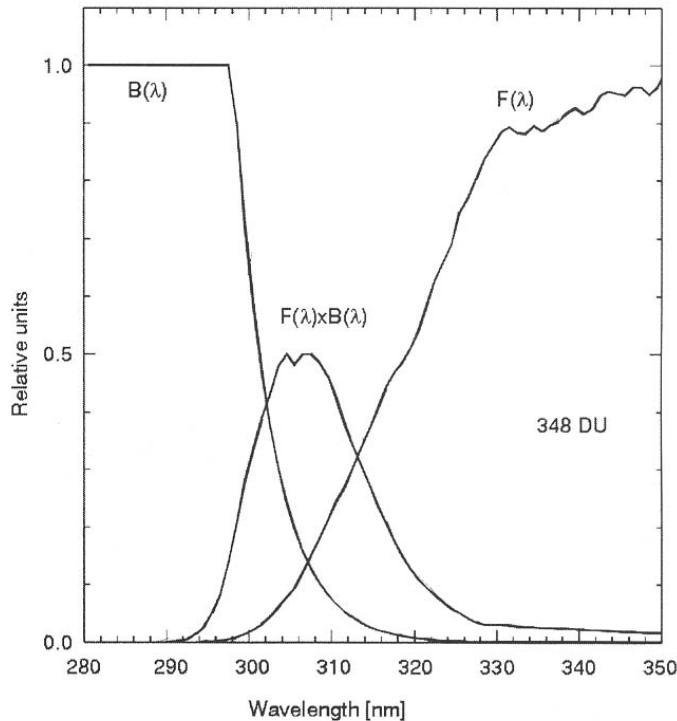


Fig. 1. Biologically active UV radiation. The overlap between the spectral irradiance  $F(\lambda)$  and the erythemal action spectrum  $B(\lambda)$  given by McKinlay and Diffey [6] shows the spectrum of biologically active radiation,  $F(\lambda)B(\lambda)$ . The area under the product function  $F(\lambda)B(\lambda)$  is the biologically active dose rate. For a total ozone column of 348 DU.

Humans are:

- strongly affected by exposure to UV-C radiation (100 to 280 nm)
- moderately affected by exposure to UV-B radiation (280 to 315 nm)
- weakly affected by exposure to UV-A radiation (315 to 400 nm)

[http://www.who.int/uv/uv\\_and\\_health/en](http://www.who.int/uv/uv_and_health/en)

← From Mandronich et al., *J. Photochemistry and Photobiology*, vol. 46, pg. 5, 1998

**The “biologically active dose rate” maximizes in the UV-B region at ~305 nm, where  $\sigma_{O_3} = 3 \times 10^{-19} \text{ cm}^2 \Rightarrow \tau(0 \text{ km}) = 2.4$  (for  $O_3$  column = 300 DU)**